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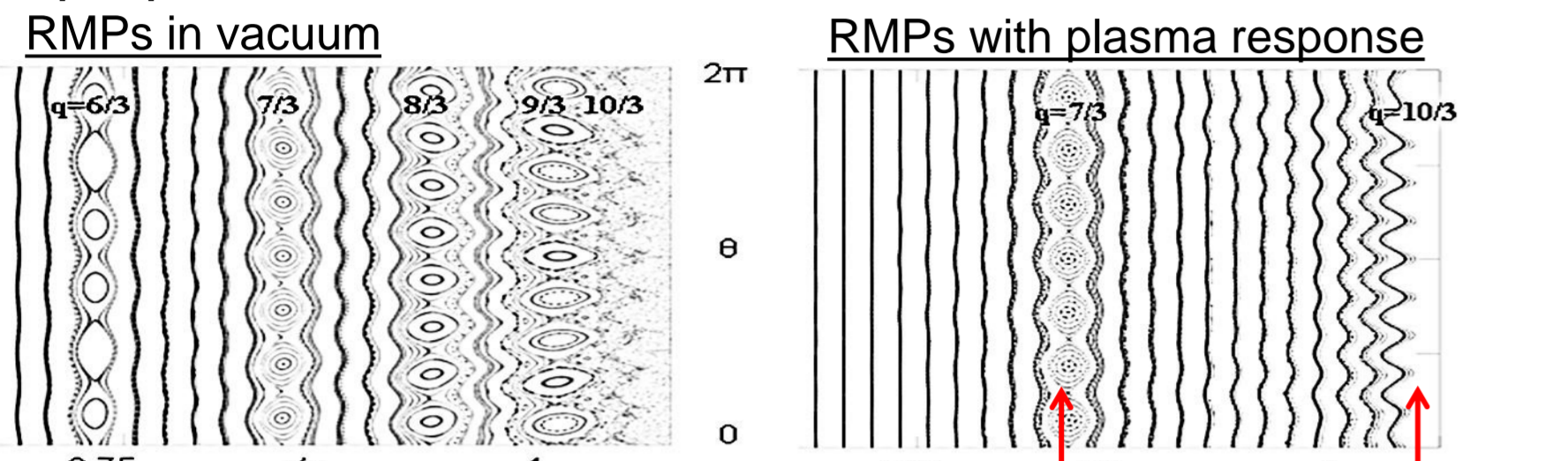
### 1/ Motivation

- Edge Localized Modes (ELMs) : Peeling-Ballooning (P-B) instabilities → needs better understanding of dynamics → mandatory control in ITER

- Promising control method: application of non-axisymmetric Resonant Magnetic Perturbations (RMPs)

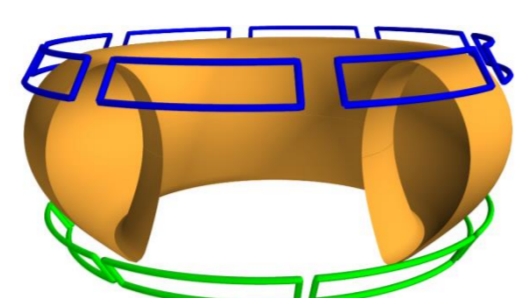
- Plasma response to RMPs not fully understood: RMPs screened or amplified by plasma flows

**Screening:** explanation: Ohm's law  $\Rightarrow \eta J_{nm} = -b_{nm}^* (V_{E,\theta} + V_{p,\theta}^e)$  perpendicular electron flow screens RMPs



Example for DIII-D case, RMP n=3 [RMHD code, Bécoulet, Orain et al, NF 2012]

"Kink response": **amplification** of stable peeling-kink modes by RMPs



AUG RMP coils [Ryan PPCF15]

**Aims:**

- better understanding of resonant (tearing) and kink responses of the plasma → effect on ELMs?
- explain features observed in experiments + predictive capabilities for ITER.

### 2/ JOEREK simulations: model and input

- JOEREK: resistive reduced MHD model with two-fluid effects: [Huysmans PPCF 2009, Orain PoP13]

Ohm's law:  $\frac{1}{R^2} \frac{\partial \psi}{\partial t} = \eta \frac{J}{R^2} - \vec{B} \cdot (\nabla \mu - \frac{e}{c} \nabla \rho)$

Mass density:  $\frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \vec{v}) + \nabla \cdot (D_{\perp} \nabla_{\perp} \rho + D_{\parallel} \nabla_{\parallel} \rho) + S_{\rho}$

Parallel momentum:  $\vec{B} \cdot (\rho \frac{d}{dt} \vec{v}_{\parallel} + \nabla_{\parallel} \rho \vec{v}_{\perp} + \nabla_{\perp} \rho \vec{v}_{\parallel}) + \nabla_{\parallel} (\rho T) + \nabla_{\perp} \cdot (\rho \vec{v}_{\perp} \otimes \vec{B} + \vec{S}_{\perp}) - \vec{v}_{\parallel} \cdot (\nabla \nabla) \vec{v}_{\parallel} = 0$

Poloidal momentum (vorticity):  $\nabla \phi \cdot \nabla \times (\rho \frac{d}{dt} \vec{v}_{\perp} + \nabla_{\perp} \rho \vec{v}_{\parallel} + \nabla_{\parallel} \rho \vec{v}_{\perp}) + \nabla_{\perp} \cdot (\rho T) + \nabla_{\parallel} \cdot (\vec{S}_{\perp} - \vec{v}_{\perp} \cdot (\nabla \nabla) \vec{v}_{\perp}) = 0$

Temperature:  $\frac{\partial (\rho T)}{\partial t} = -\vec{v}_{\perp} \cdot \nabla_{\perp} (\rho T) - \gamma \rho T \nabla_{\parallel} \cdot (\vec{v}_{\perp} + \vec{v}_{\parallel} \nabla_{\parallel}) + \nabla_{\parallel} \cdot (K_{\perp} \nabla_{\perp} T + K_{\parallel} \nabla_{\parallel} T) + (1-\gamma) S_T + 0.5 T^2 S_p$

Flows included in the model:

- ExB + diamagnetic rotation:

$$\vec{v}_i = -R^2 \nabla \mu \times \nabla \phi - \tau_{IC} \frac{R^2}{\rho} \nabla p \times \nabla \phi + V_{\parallel, i} \vec{B}$$

$$\vec{v}_E = \vec{E} \times \vec{B} \quad \text{diamagnetic}$$

- Source of toroidal rotation  $S_{\phi}$

- Neoclassical poloidal friction:

$$\nabla_{\perp} \cdot \Pi_{i, neo} \approx \mu_{i, neo} \rho (B^2 / B_0^2) (\vec{v}_{\perp, i} - \vec{v}_{\perp, neo}) \cdot \vec{e}_{\theta}$$

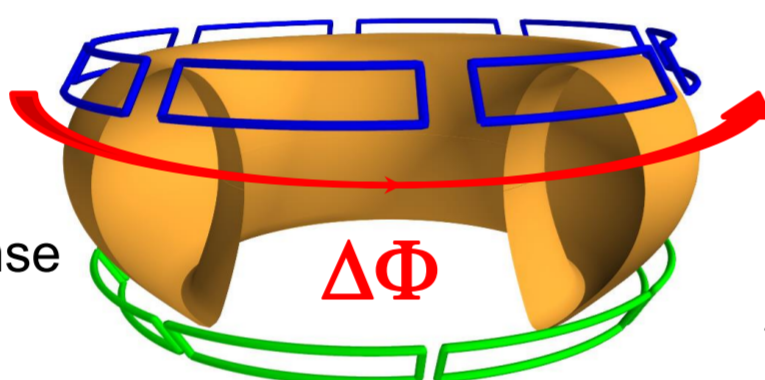
$$V_{\theta, i} \rightarrow V_{\theta, neo} = -k_{i, neo} \tau_{IC} (\nabla_{\perp} \cdot \nabla_{\perp} T) / B_0$$

- Main numerical limitation: resistivity  $\eta \approx 10-100 \times \eta_{Spitzer}$

- Input: ASDEX Upgrade discharges (#31128, #30835) ELM mitigation with n=2 RMPs

- RMP spectrum applied as boundary condition

- 1/ RMP field calculated in vacuum (VACFIELD)
- 2/ Applied at the boundary of JOEREK domain and increased in 1000 t<sub>k</sub> → penetration takes into account plasma response



Rotation of differential phase  $\Delta\Phi$  between upper and lower coil currents

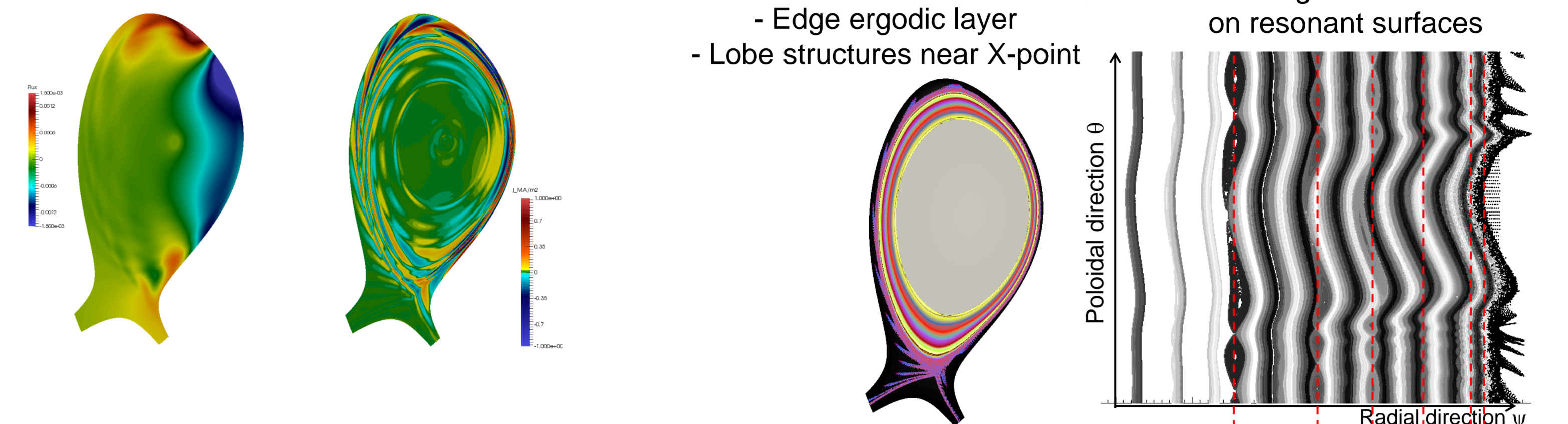
- change applied RMP spectrum
- change plasma response to RMP

- Simulation: 1/ Axisymmetric equilibrium from equilibrium reconstruction with CLISTE [Dunne NF2012] 2/ add n≠0 modes → ELM growth without/with RMPs

### 3/ Plasma response to RMPs

- General features of plasma response:**

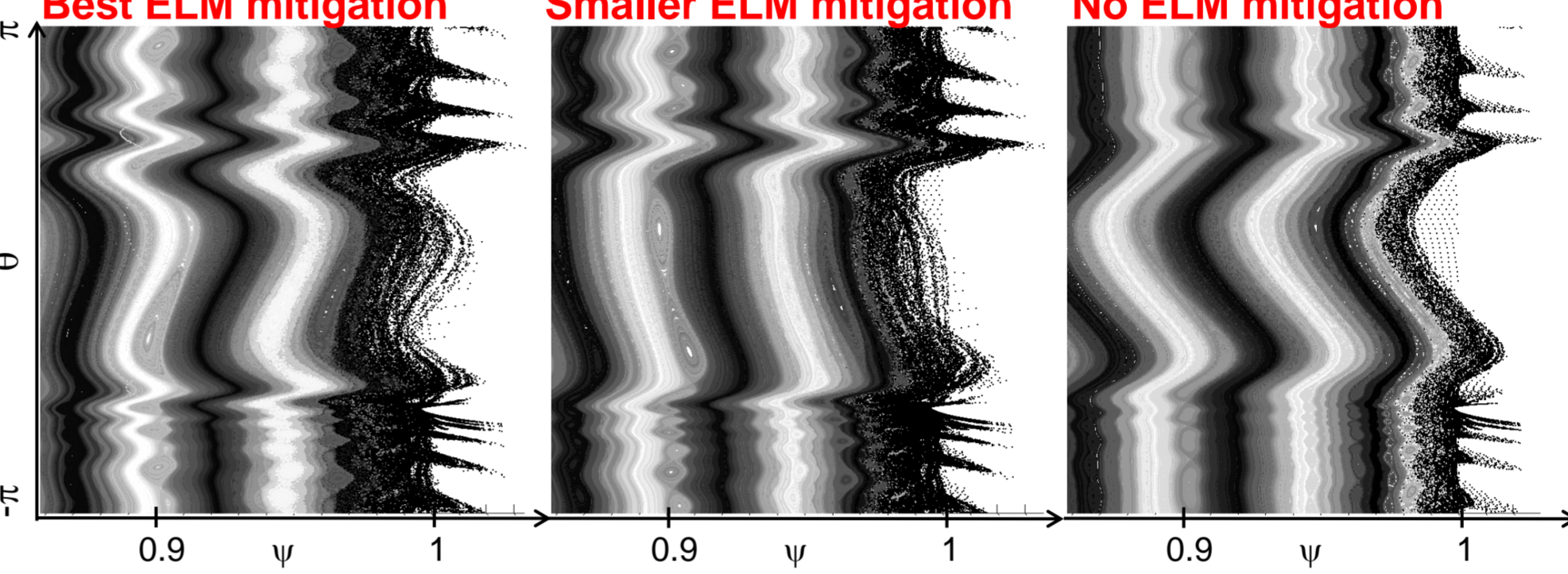
- Penetration of n=2 magnetic flux perturbation
- Response currents on resonant surfaces
- Magnetic island chains induced on resonant surfaces where  $V_{\perp, e}$  small + ergodic layer for  $\psi \geq 0.98$



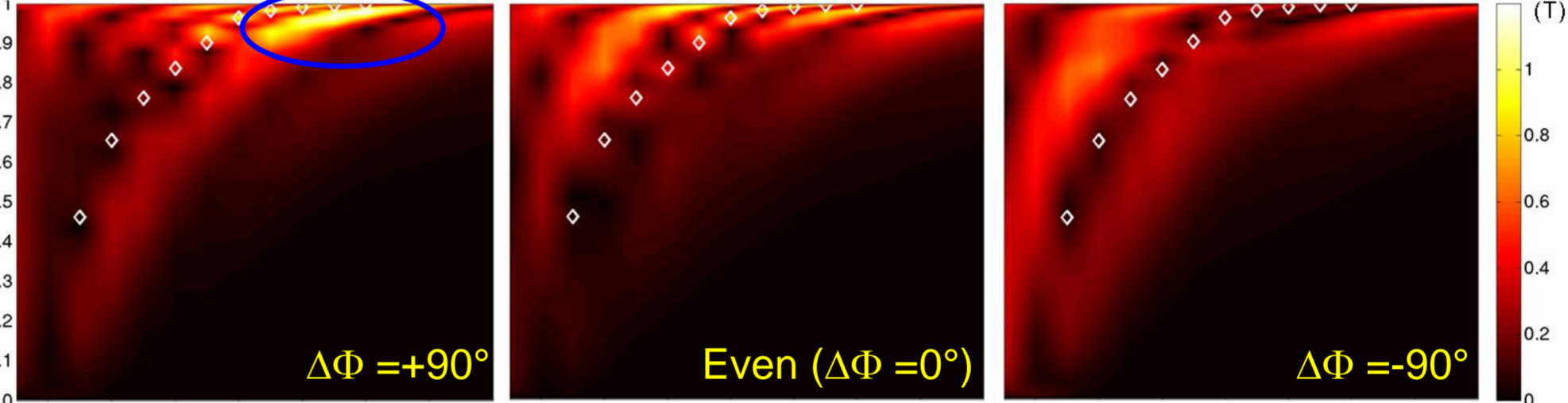
- Influence of applied RMP spectrum :**

- Larger ergodic layer in case  $\Delta\Phi = +90^\circ$  + larger field line kinking near X-point

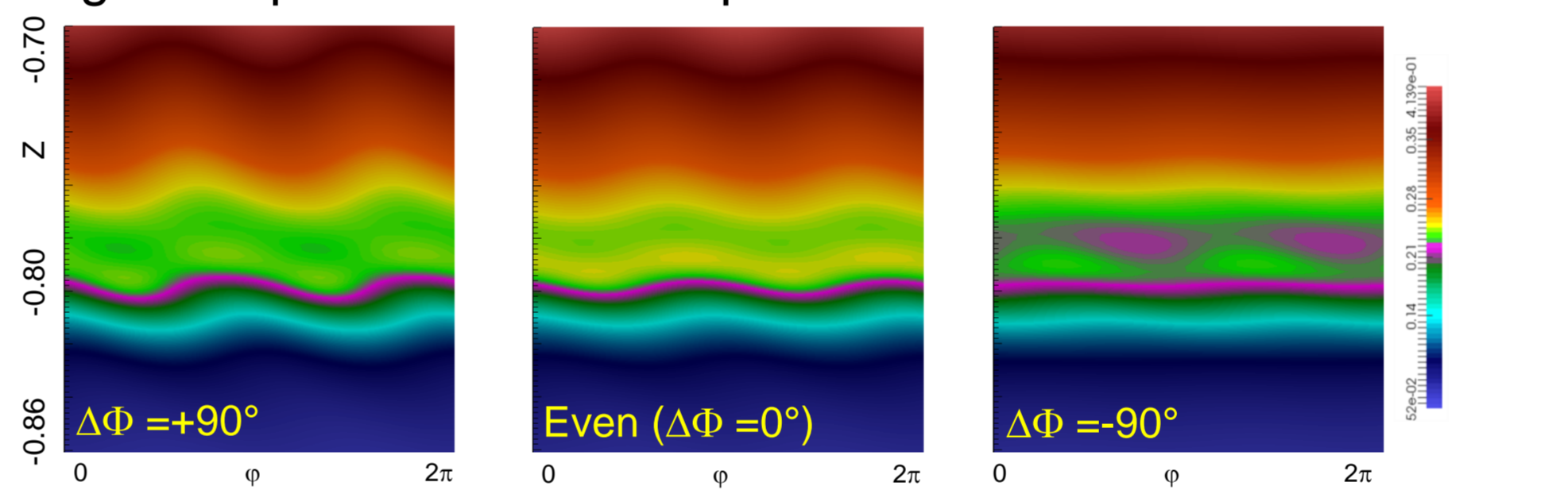
- $\Delta\Phi = +90^\circ$  Best ELM mitigation
- Even ( $\Delta\Phi = 0^\circ$ ) Smaller ELM mitigation
- $\Delta\Phi = -90^\circ$  No ELM mitigation



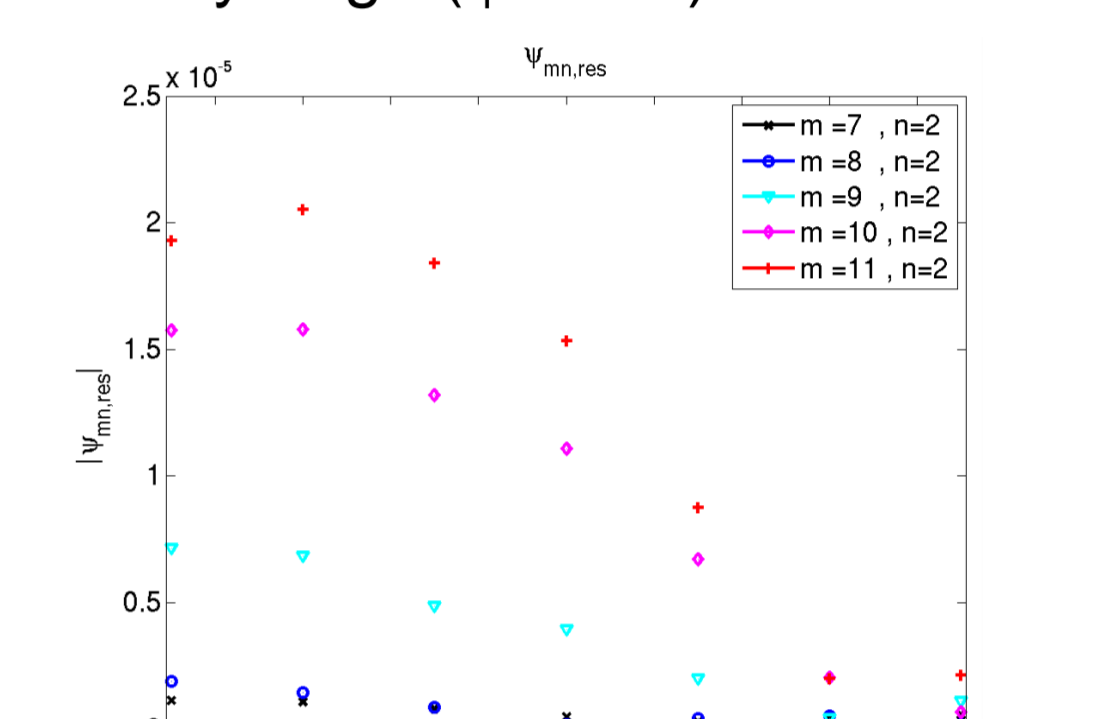
- Largest edge kink response in case  $\Delta\Phi = +90^\circ$   $\delta Br$  (n=2):



- Largest displacement near X-point in case  $\Delta\Phi = +90^\circ$ :



- Screening in pedestal in all cases.  $\Delta\Phi = +60$  to  $+90^\circ$ : most resonant at very edge ( $\psi > 0.97$ )



- Case  $\Delta\Phi = +90^\circ$ : strongest ELM mitigation in experiments

- In modeling: Coupling of  $m > nq$  'kink mode' with  $m$  resonant component → amplification of resonant response

[Orain et al, NF 2016]

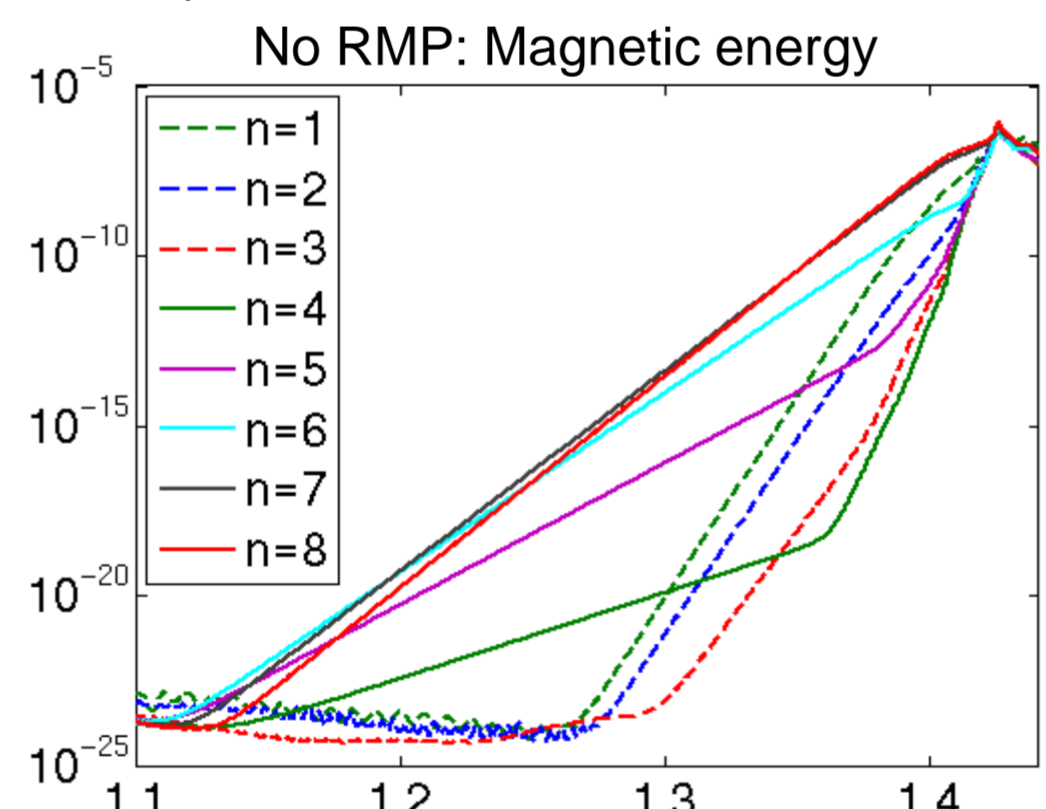
### 4/ ELM/RMP interaction

- Modeling for #31128 (multi-harmonic n=1-8) of ELM without and with RMP for  $\Delta\Phi = +90/-90^\circ$ :

- Start from P-B unstable profiles (n=8 most P-B unstable)

- ELM without RMP:

First linear growth of modes, then non-linear mode coupling, mode saturation and ELM crash:

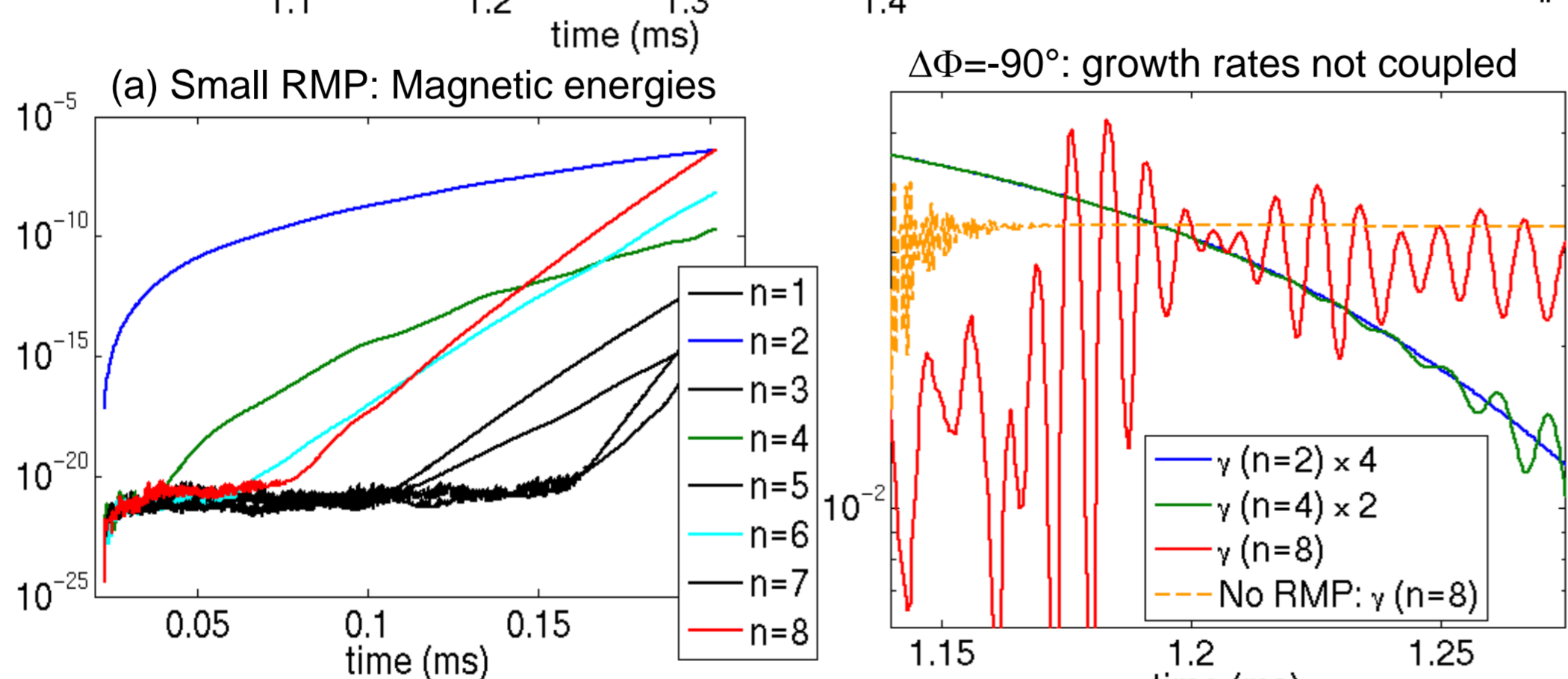


- ELM with RMP:

- 2 opposite behaviours:

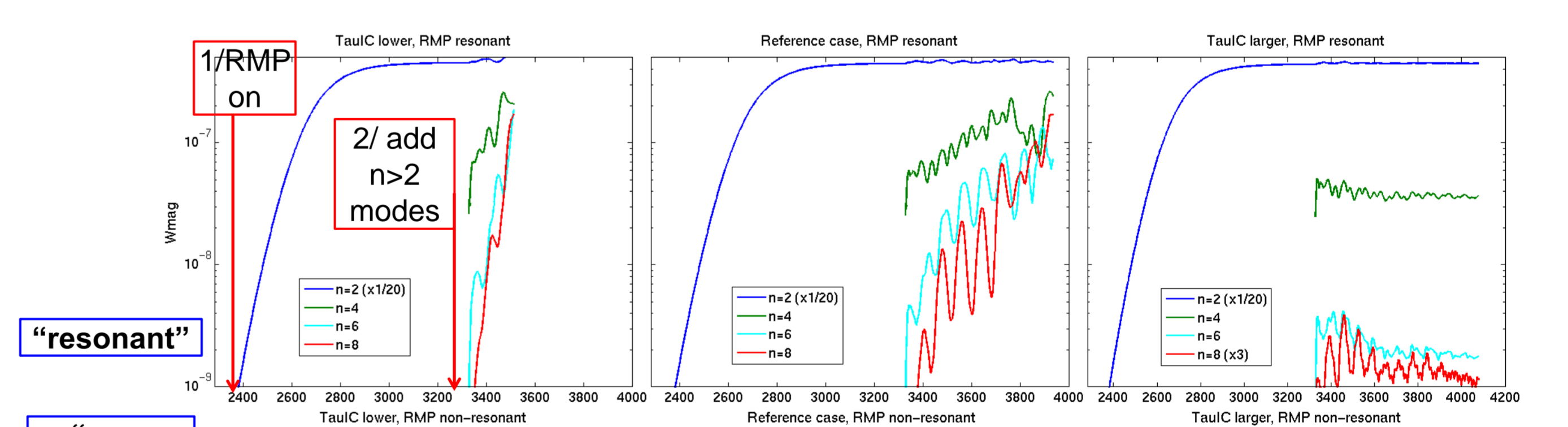
- (a) for small "penetrated" RMP amplitude:

n=8 and n=6 not coupled to RMPs → modes grow as unmitigated ELM, with only reduced growth rate  $\gamma$



- (b) for large penetrated" RMP amplitude or large viscosity: even modes driven by n=2 RMP → full stabilization of ELM

- Scan in diamagnetic rotation: → When increasing perpendicular diamagnetic rotation  $V_{dia}$ : bifurcation from unstable ELM to stabilized ELM (ELM suppression?) via partially stabilized ELM (ELM mitigation?) → Threshold depends on resonant spectrum (kink amplification)



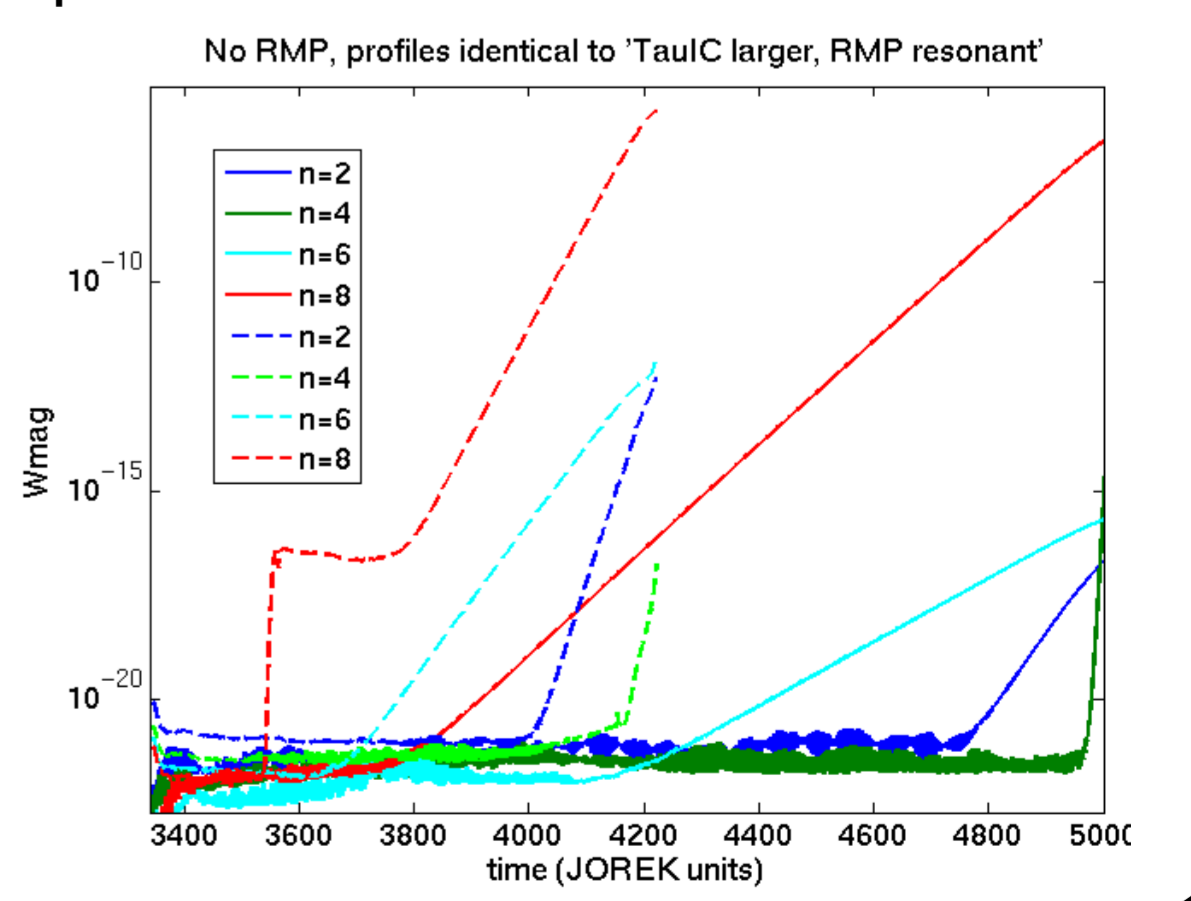
- Other important parameters:

- resistivity  $\eta$  ++: destabilizing since makes ELMs more unstable
- viscosity  $\nu$  ++: stabilizing since increases RMP penetration

ELM mitigation/suppression = f( $\nu, \eta, V_{dia}$ ) + RMP amplitude + other parameters?

- Simulation with pressure profile degraded (similar to RMP effect) but without RMP:

→ Reduction of pressure gradient by RMP is NOT responsible for ELM suppression → needs presence of mode coupling to get QH-mode-like behaviour.



### Conclusions

- Plasma response to n=2 RMPs and ELM/RMP interaction modeled in JOEREK using AUG experimental input.

- Plasma response to RMPs:

Best ELM mitigation related to strongest kink response → mechanism: coupling between  $m > nq$  kink component with m resonant component → amplification of resonant perturbation. → Consistent with MARS-F [Ryan et al, PPCF2015] and NEMEC [Strumberger 2014] modeling.

- ELM/RMP interaction:

Depending on plasma response, either coupling of high-n modes with RMPs or barely affected development of ELM. → ELM mitigation/suppression NOT due to reduction of pressure gradient by RMPs → ELM mitigation possibly induced by strong coupling of even modes with n=2 RMP + damping of odd modes. → changing ELMs into saturated peeling-kink modes?

- Perspectives: Current and future works:

- Further understanding of ELM mitigation or suppression by RMPs with quantitative comparisons to experiments
- improved model to observe density pumpout in simulations
- RMP coils directly in JOEREK-STARWALL [Hoelzl et al, JCP 2012]
- Modeling of ELM cycles with/without RMPs