

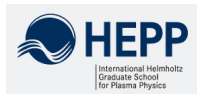
# Non-linear Simulation of Edge Localized Modes in ASDEX Upgrade

**Alexander Lessig,**

Matthias Hölzl, François Orain, Sibylle Günter, Marina Becoulet, Guido Huysmans, Stanislas Pamela, Mike Dunne, the ASDEX Upgrade Team



*EUROfusion*



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## 2 JOREK: Nonlinear MHD

- Overview

- Reduced MHD

- Numerics

## 3 Modelling of ELM Dynamics

- Full Crash

- Resolution Scan

- ELM Losses/Divertor Heat Load

- Filament Dynamics

## 4 Summary and Outlook

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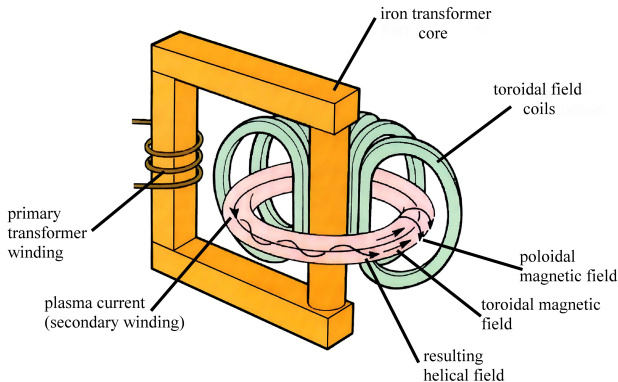
Full Crash

Resolution Scan

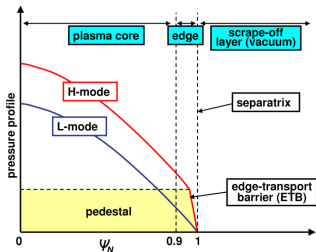
ELM Losses/Divertor Heat Load

Filament Dynamics

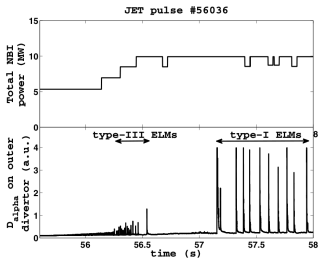
## 4 Summary and Outlook



- ▶ magnetic field in toroidal direction created by toroidal field coils
  - ▶ inductively driven plasma current creates poloidal magnetic field
- plasma confinement by helically winding field lines



- ▶ relaxation-like oscillatory instability at the boundary of H-mode plasmas (pedestal collapse)
- ▶ driven by large pressure gradients and current densities
- ▶ eject particles and energy on very short time-scales
- ▶ plasma particle and energy content decreased by up to 10%



- ▶ **high heat loads on plasma-facing components (in particular divertor targets)**

→ comprehensive theoretical understanding necessary for control and prediction of ELM properties

S. Pamela. Ph.D. thesis, University of Provence (2010)

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▶ **3D Non-linear MHD in X-point tokamak geometry**

G. Huysmans and O. Czarny. *Nucl Fusion*, 47, 659 (2007)

- ▶ European project
- ▶ multi-purpose non-linear MHD code

▶ **Edge Localized Modes in**

- ASDEX Upgrade M. Hölzl, S. Günter, et al. *Phys Plasmas*, 19, 082505 (2012)
- MAST S. J. P. Pamela, G. T. A. Huysmans, et al. *PPCF*, 55, 095001 (2013)
- JET S. J. P. Pamela, G. T. A. Huysmans, et al. *PPCF*, 53, 054014 (2011)
- ITER G. Huysmans and A. Loarte. *J Nucl Mater*, 438, s57 (2013)

▶ **Resonant Magnetic Perturbations** M. Bécoulet, F. Orain, et al. *Phys Rev Lett*, 113, 115001 (2014)

▶ **Pellet ELM Triggering** G. Huysmans, S. Pamela, et al. *23rd IAEA, THS/7-1* (2010)

▶ **Tearing Modes** J. Pratt and E. Westerhof. *54th APS* (2012)

▶ **Disruptions** A. Fil, E. Nardon, et al. *Phys Plasmas*, 22, 062509 (2015)

▶ ...



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- ▶ **Starting point: resistive MHD equations** including particle and heat sources/diffusion

$$\partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = \nabla \cdot (D_{\perp} \nabla_{\perp} \rho) + S_{\rho}$$

$$\rho (\partial_t + \mathbf{v} \cdot \nabla) \mathbf{v} = -\nabla p + \mathbf{j} \times \mathbf{B} + \mu \Delta \mathbf{v}$$

$$\partial_t p + \mathbf{v} \cdot \nabla p + \gamma p \nabla \cdot \mathbf{v} = \nabla \cdot (\kappa_{\parallel} \nabla_{\parallel} T + \kappa_{\perp} \nabla_{\perp} T) + S_T$$

$$\mathbf{E} + \mathbf{v} \times \mathbf{B} = \eta \mathbf{j}$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

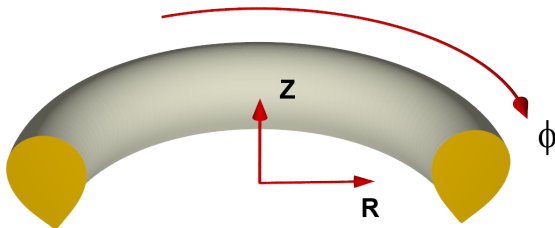
$$\nabla \times \mathbf{B} = \mu_0 \mathbf{j}$$

$$\nabla \cdot \mathbf{B} = 0$$

- ▶ Variables:  $\rho$ ,  $\mathbf{v}$ ,  $p$ ,  $\mathbf{B}$  with  $p = R_s \rho T$

- ▷ Definition of magnetic and velocity vector fields:  $\mathbf{B} := \underbrace{\frac{F_0}{R} \mathbf{e}_\phi}_{\mathbf{B}_\phi} + \underbrace{\frac{1}{R} \nabla \Psi \times \mathbf{e}_\phi}_{\mathbf{B}_{\text{pol}}}$

and  $\mathbf{v} := \underbrace{-R \nabla u \times \mathbf{e}_\phi}_{\mathbf{v}_\perp} + \underbrace{v_{\parallel} \mathbf{B}}_{\mathbf{v}_{\parallel}}$  with  $|\mathbf{B}_\phi| \gg |\mathbf{B}_{\text{pol}}|$



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$$\partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = \nabla \cdot (\mathbf{D}_\perp \nabla_\perp \rho) + S_\rho$$

$$\mathbf{e}_\phi \cdot \nabla \times R^2 \left\{ \rho (\partial_t + \mathbf{v} \cdot \nabla) \mathbf{v} = -\nabla p + \mathbf{j} \times \mathbf{B} + \mu \Delta \mathbf{v} \right\}$$

$$\mathbf{B} \cdot \left\{ \rho (\partial_t + \mathbf{v} \cdot \nabla) \mathbf{v} = -\nabla p + \mathbf{j} \times \mathbf{B} + \mu \Delta \mathbf{v} \right\}$$

$$\partial_t p + \mathbf{v} \cdot \nabla p + \gamma p \nabla \cdot \mathbf{v} = \nabla \cdot (\kappa_{\parallel} \nabla_{\parallel} T + \kappa_{\perp} \nabla_{\perp} T) + S_T$$

$$\partial_t \Psi = R[\Psi, u] + \frac{\eta}{\mu_0} \mathbf{j} - F_0 \partial_\phi u$$

▶ Definitions:  $\omega := \nabla_\perp^2 u$  and  $\mathbf{j} := \Delta^* \Psi$

▶ Variables:  $\rho, \omega, u, v_{\parallel}, p, \Psi, \mathbf{j}$  with  $p = R_s \rho T$

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- ▶ **toroidal direction:** Fourier decomposition
- ▶ **poloidal plane:** 2D Bezier finite elements ( $C^1$ , iso-parametric)

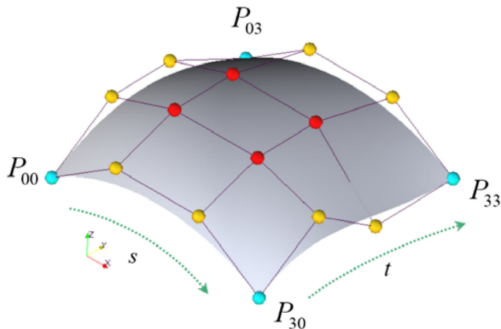
$$\mathbf{P}(s, t) = \sum_{i=0}^3 \sum_{j=0}^3 \mathbf{P}_{ij} B_i^3(s) B_j^3(t)$$

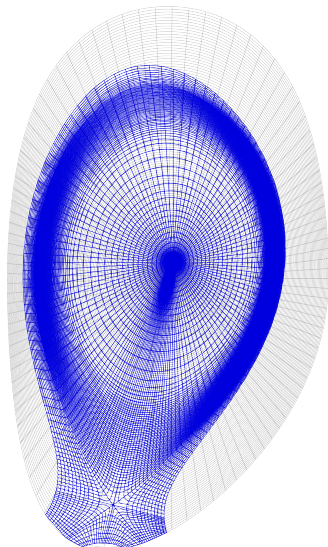
$$B_0^3(s) = (1 - s)^3$$

$$B_1^3(s) = 3(1 - s)^2 s$$

$$B_2^3(s) = 3(1 - s) s^2$$

$$B_3^3(s) = s^3$$





- ▶ Construction of polar grid
- ▶ Solution of Grad-Shafranov equation on polar grid using input profiles ( $F_0$ ,  $\Psi_{\text{bnd}}$ , profiles for  $T$ ,  $\rho$ ,  $FF'$ )
- ▶ Construction of flux-aligned X-point grid
- ▶ ideal Wall and Bohm boundary conditions
- ▶ Fully implicit time integration (GMRES, Physics-based preconditioner, hybrid openMP/MPI parallelization)
- ▶ Postprocessing

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**Full Crash**

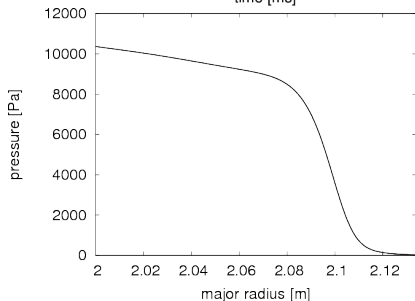
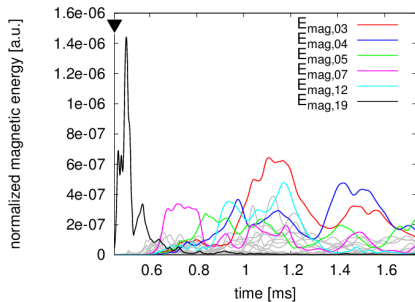
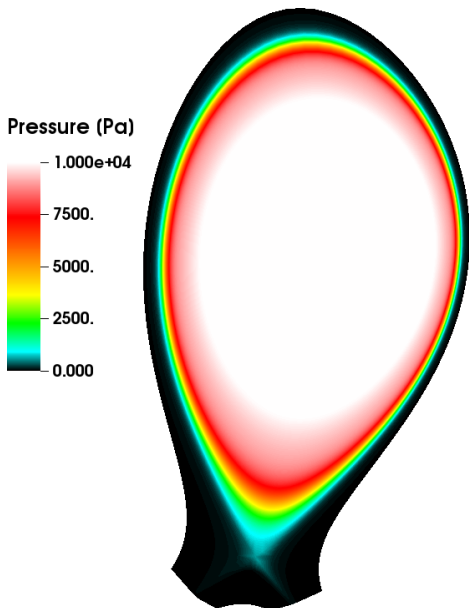
Resolution Scan

ELM Losses/Divertor Heat Load

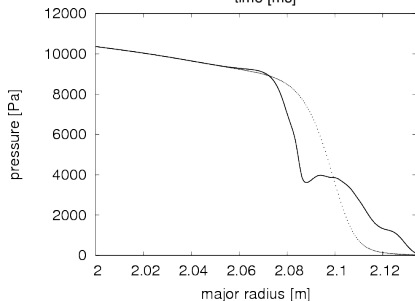
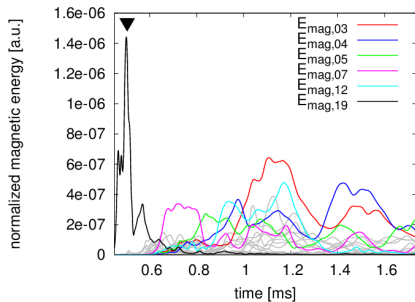
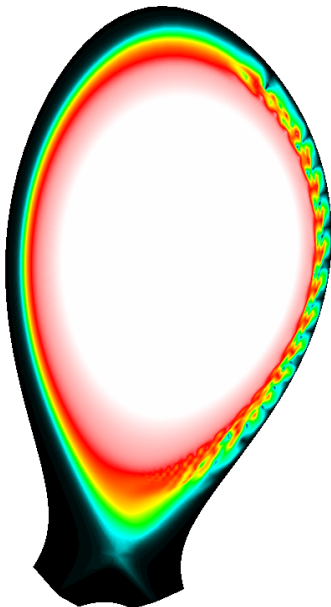
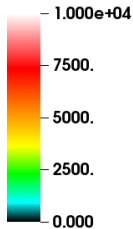
Filament Dynamics

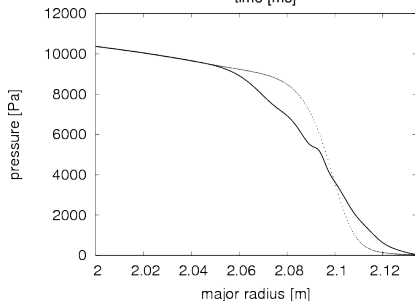
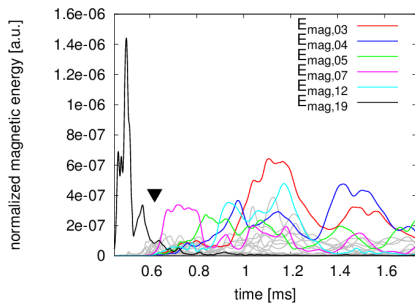
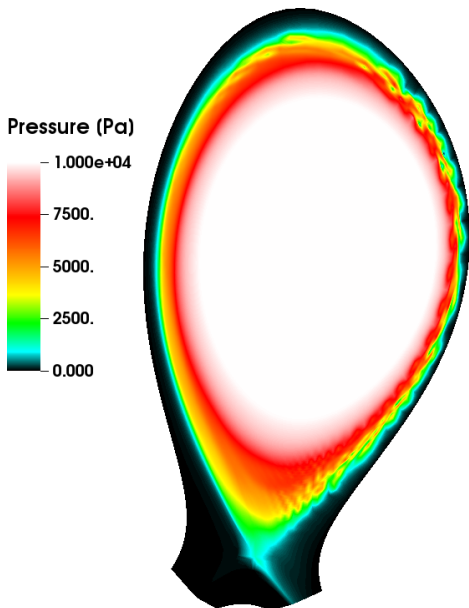
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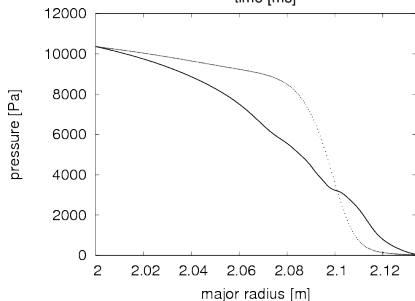
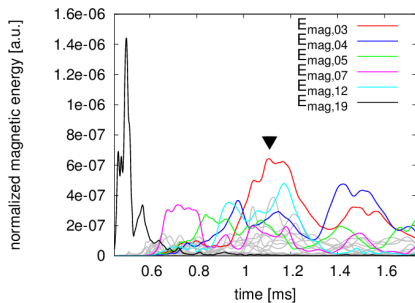
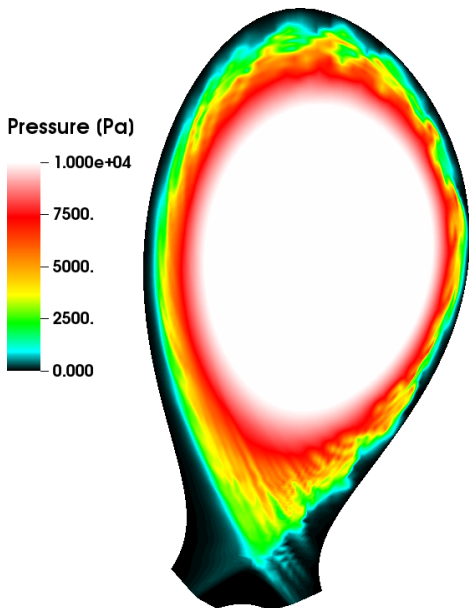
- ▶ ASDEX Upgrade discharge #29342@4.25s CLISTE reconstruction by Mike Dunne
- ▶ toroidal modes  $n = 0 \dots 22$  included
- ▶ core resistivity in ASDEX Upgrade:  $\sim 1 \cdot 10^{-8} \Omega\text{m}$
- ▶ core resistivity in simulation:  $2.5 \cdot 10^{-7} \Omega\text{m}$
- ▶ about 500,000 cpu-hours on Helios



Pressure (Pa)







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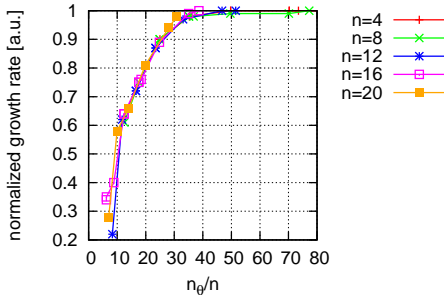
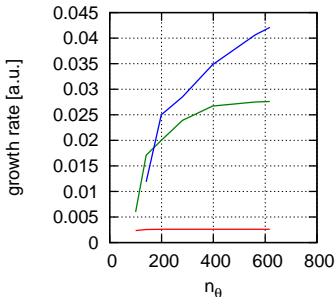
Full Crash

**Resolution Scan**

ELM Losses/Divertor Heat Load

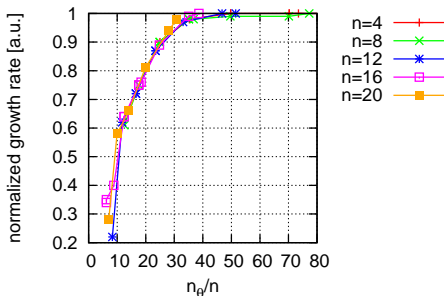
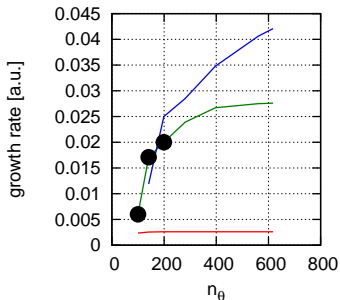
Filament Dynamics

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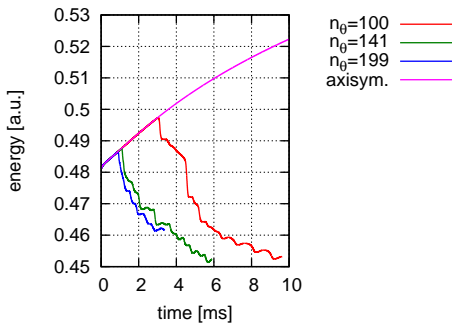


- ▶ ASDEX Upgrade discharge #31128
- ▶ core resistivity in simulation:  $2.5 \cdot 10^{-7} \Omega\text{m}$  (AUG:  $\sim 1 \cdot 10^{-8} \Omega\text{m}$ )
- ▶ high- $n$  components require high grid resolution ( $m = q \cdot n$ ) → **linearly not fully resolved**





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- ▶ high-n components require high grid resolution ( $m = q \cdot n$ ) → **linearly not fully resolved**
- ▶ losses converged → **non-linearly well resolved**

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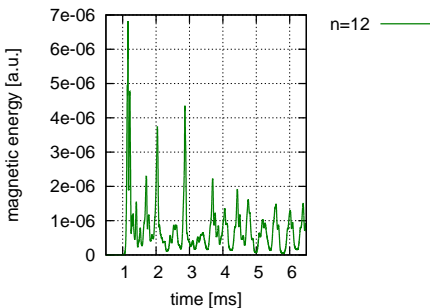
**ELM Losses/Divertor Heat Load**

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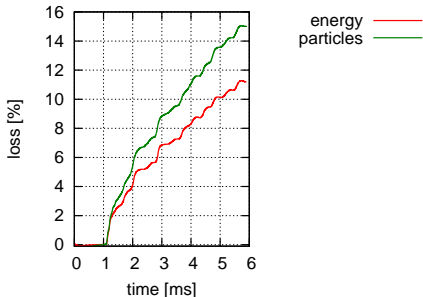
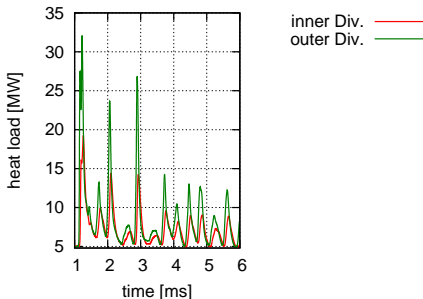
## Modelling

## Losses/Divertor Heat Load



- ▶ main ELM phase for  $t \leq 3\text{ms}$
- ▶ ballooning turbulence for  $t > 3\text{ms}$
- ▶ particle/energy losses of  $\sim 10\%$
- ▶ asymmetry in divertor heat load
- ▶ diamagnetic drift influences heat load asymmetry and ballooning turbulence

F. Orain, M. Becoulet, et al. *Phys Rev Lett*, 114, 035001 (2015)



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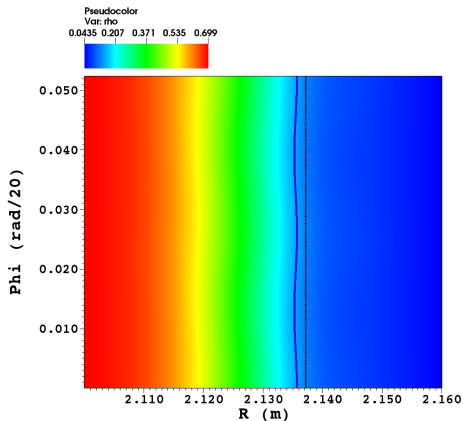
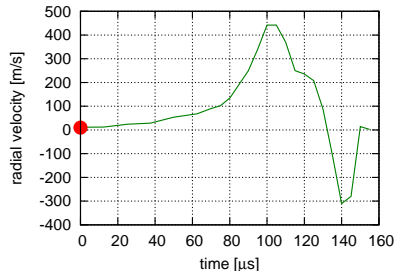
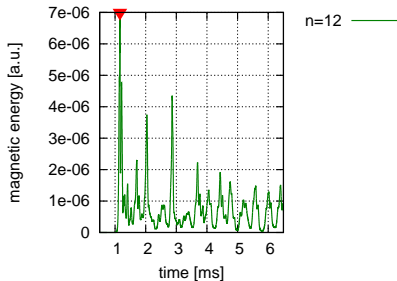
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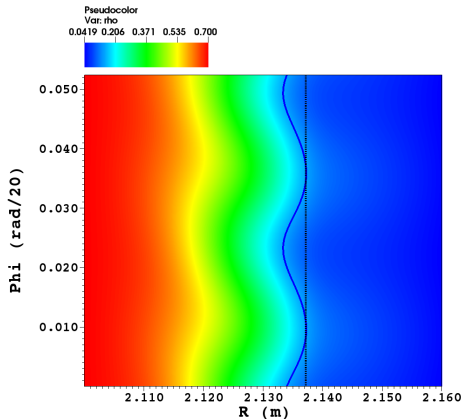
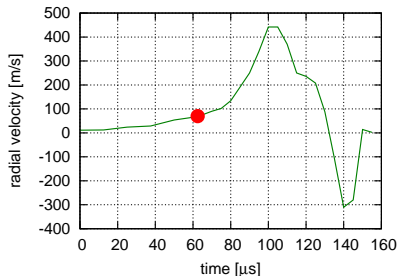
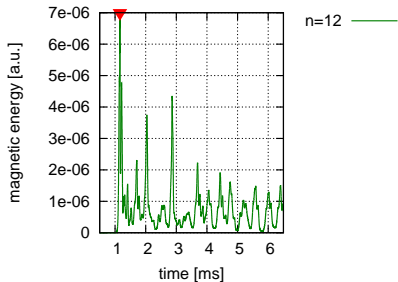
ELM Losses/Divertor Heat Load

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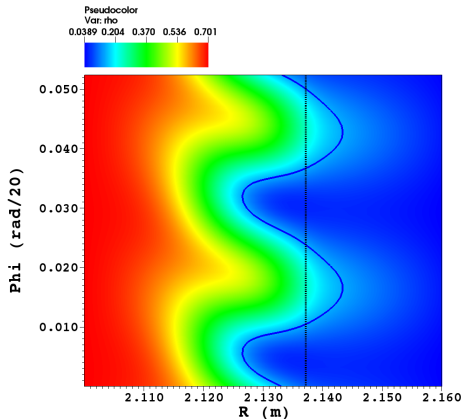
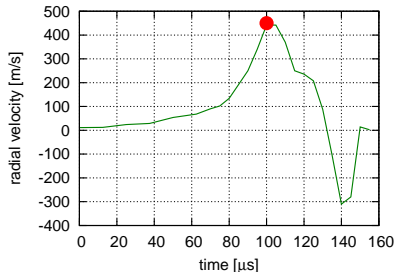
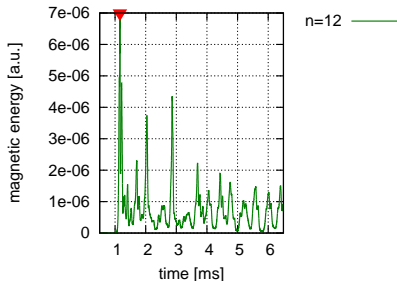
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- ▶ radial size of filament:  $\sim 4$ cm
- ▶ poloidal size of filament:  $\sim 13.5$ cm

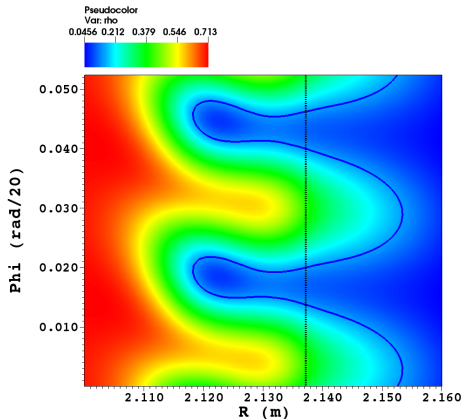
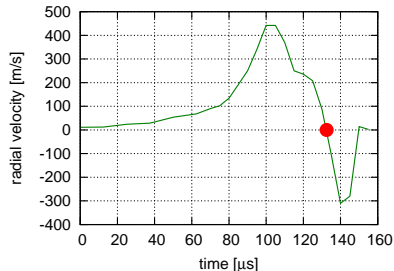
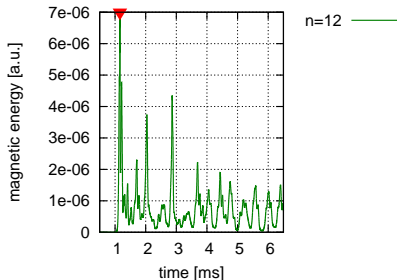


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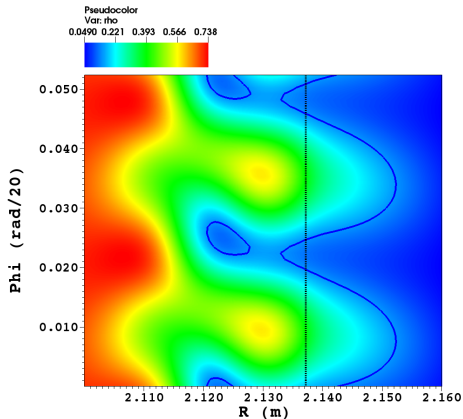
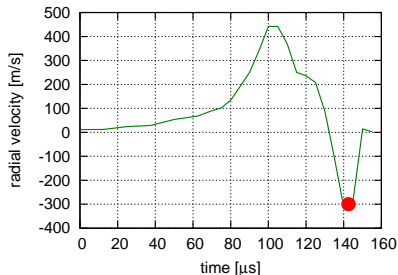
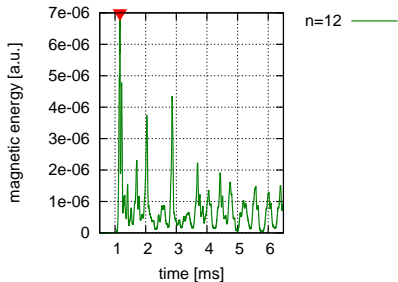


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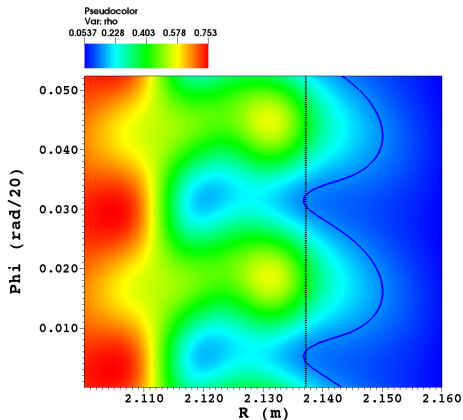
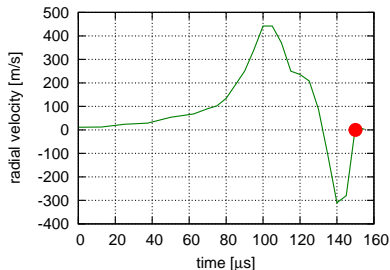
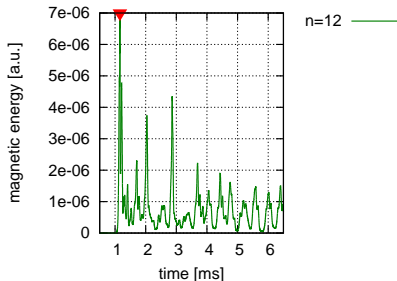




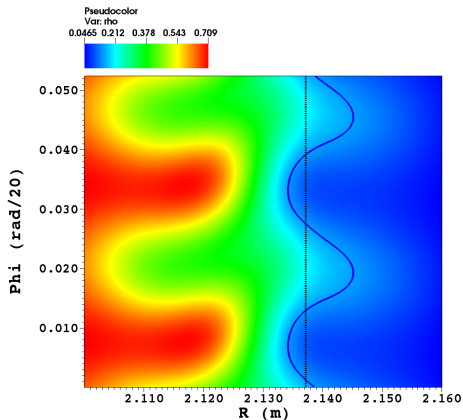
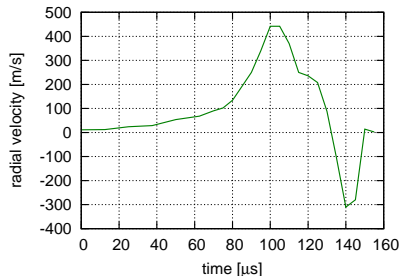
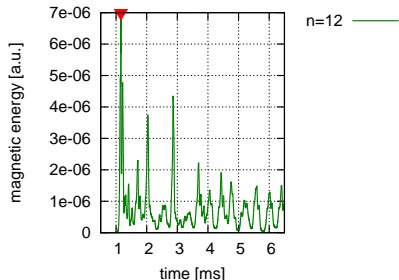
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## Summary

- ▶ high-n components require high grid resolution → computational limits in linear phase/non-linearly well resolved
- ▶ ELM losses comparable to experimental observations for type-I ELMs
- ▶ evolution of toroidal Fourier spectrum and pedestal in qualitative agreement with experimental observations

H. Zohm. *PPCF*, 38, 105 (1996)

R. P. Wenninger, H. Reimerdes, O. Sauter, and H. Zohm. *Nucl Fusion*, 53, 113004 (2013)

P. Schneider, E. Wolfrum, et al. *PPCF*, 56, 025011 (2014)

## Outlook

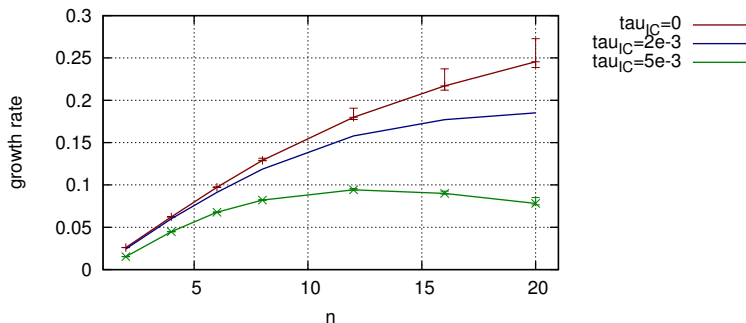
- ▶ include diamagnetic drifts (act stabilizing on high mode numbers)
- ▶ include toroidal rotation and neoclassical viscosity
- ▶ identify different ELM types in our simulations as observed in experiments
- ▶ further comparison to experimental observations, e.g. pedestal evolution and heat deposition patterns

# References

- M. Bécoulet, F. Orain, et al.** *Phys Rev Lett*, 113, 115001 (2014).
- O. Czarny and G. Huysmans.** *J Comput Phys*, 227, 7423 (2008).
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- ▶ diamagnetic drift terms stabilize high- $n$  components
- ▶ intermediate mode numbers become dominant