

Steady state heat transport in magnetized plasmas with magnetic islands and local stochastic fields

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Motivation	Model	Magnetic islands	Stochastic layers	Summary	References

Motivation

Model

Magnetic islands

Stochastic layers

Summary

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Motivation

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Motivation: Magnetic islands

- magnetic reconnection (example: 3/2-island)
- field lines wander around island surfaces
- parallel transport ($\chi_{||}$): fast, long distance
- perpendicular transport (χ_{\perp}): slow, short distance
- ▶ parallel + perp. transport ⇒ temperature flattening
- ▶ scale island width w_c : parallel \approx perpendicular transport



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Motivation: Neoclassical tearing mode

- bootstrap current: toroidal current driven by $\nabla_r p$
- perturbed by island temperature flattening
 - effective lack current in the island o-point region
 - acts as a driving term for further island growth
 - ► ⇒ neoclassical tearing mode (NTM)
- NTM stability strongly depends on temp. distribution
- exact heat flux computations are important
- ► analytical theory is limited to the cases $w/w_c \rightarrow 0$ and $w/w_c \rightarrow \infty$, see Fitzpatrick (1995)
- AIM: numerical computations for realistic parameters

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Motivation: Stochastic layers

overlapping magnetic islands destroy flux surfaces



- field lines move through stochastic layer in a seemingly random way (increases radial heat transport)
- AIM: determine radial heat conductivity of highly stochastic layers and compare to analytical predictions

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Model

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Model: Geometry

- equilibrium: circular cross section, large aspect ratio
- q-profile with 0.9 in the center and 4.0 at the edge



- coordinates: helical, unsheared, rational helicity q_c (adapted to problem)
- grid points in radial and poloidal direction
- Fourier expansion in toroidal direction

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Model: Heat diffusion equation

- Solve steady state heat diffusion equation $\nabla \vec{q} = P$
 - P: power source (heating)
 - $\vec{q} = -n\chi_{||}\nabla_{||}T n\chi_{\perp}\nabla_{\perp}T$: heat flux density
 - *n*: particle density
 - $\chi_{||}/\chi_{\perp}$ typically between 10⁷ and 10¹¹
- ► common numerical schemes: error $\propto \chi_{||}/\chi_{\perp} \Rightarrow$ exact coordinate alignment
- virtually impossible for time-dependent problems
- new scheme developed by Günter et al. (2005)
 - conserves self-adjointness of parallel transport operator
 - temperature and heat flux grids shifted against each other
 - Fourier cut-off is performed at a certain heat flux order

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Model: New numerical scheme

- deviations from analytical solution for non-trivial test case
- small numerical errors
- independent from $\chi_{||}/\chi_{\perp}$ up to 10¹³



code benchmarks were also performed

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Magnetic island results

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Magnetic islands: Radial heat diffusivity χ_r

- effective radial heat diffusivity χ_r
 - increased in the island region
 - maximum at the resonant surface
- ► 4/3-island with w = 0.068a:



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Magnetic islands: Scaling of χ_r

• χ_r at the island resonant surface



• two different regimes:
$$\chi_r \propto \left(\frac{w}{w_c}\right)^4$$
 resp. $\chi_r \propto \left(\frac{w}{w_c}\right)^2$

• depends on w/w_c only

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Magnetic islands: Heat flux

- ▶ Total heat flux in the island region (4/3-island, w = 0.068a)
- $w/w_c = 3.4 \Rightarrow$ largely flattened island



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Magnetic islands: Heat flux components

- Same case
- Parallel (red) and perpendicular (blue) heat flux



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NTM stability: Comparison to analytical limits

- NTMs destabilized by seed island temperature pert.
- Neoclassical contribution to island growth rate:



Numerical results match analytical limits

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NTM stability: Comparison to analytical matching

Fitzpatrick performed matching of the analytical limits



- underestimates island growth rate significantly
- can make the difference between stable and unstable

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Results for highly stochastic layers

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Stochastic layers: Flux surface destruction

- ► flux surfaces destroyed for $s = \frac{(w_1+w_2)/2}{|r_{res,1}-r_{res,2}|} \gtrsim 1$
- test case:
 - ▶ 5 islands: 24/23, 25/24, 26/25, 27/26, 28/27
 - total stochasticity $s = 48.5 \gg 1$



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Stochastic layer: Radial heat diffusivity χ_r



radial heat diffusivity \(\chi_r\) strongly increased in the region of the stochastic layer

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Stochastic layers: Analytical theories

- ► analytical theories for heat transport across highly stochastic layers (s ≫ 1): review by Liewer (1985)
- Rechester-Rosenbluth regime
 - Iow collisionality
 - electrons basically follow the stochastic field lines
 - transport dominated by field line diffusion
- Kadomtsev-Pogutse regime
 - medium collisionality
 - increased importance of electron diffusion
- fluid regime
 - high collisionality
 - transport dominated by electron diffusion

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Stochastic layer: Scaling of χ_r

• χ_r in the center of the ergodic layer:



- the three analytically predicted regimes can be observed
- ranges of validity do not coincide

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Summary

- implemented code for heat diffusion computations
- demonstrated computations with unaligned coordinates
- radial heat diffusivity χ_r for islands
- [dw/dt]_{bs} for neoclassical tearing modes; widely used analytical matching underestimates island growth!
- χ_r at highly stochastic layers with realistic parameters
 - found the analytically predicted regimes
 - different ranges of validity
- experience from this work will be used to implement a nonlinear MHD code for plasma edge examinations



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Island code benchmark with TM1

> 1/1 magnetic perturbation with w = 0.118a

• at
$$\chi_{||}/\chi_{\perp} = 10^{11}$$
: $w/w_c = 13.8$



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Ergodic code benchmark with TM1

- 4/3 and 3/2 magnetic perturbations
- ► s = 1.4
- w/wc between 0.34 and 3.4
- Relative differences for representative modes:



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Analytical test case

$$ec{B} =
abla \Psi imes \hat{\mathbf{e}}_{\phi},$$
 (1)

$$\Psi(r,\theta,\phi) = \Psi_0(r) + \Psi_1(r,\theta,\phi), \tag{2}$$

$$\Psi_0(r) = \Psi_0 \left(\frac{r}{a}\right)^2 \left(\frac{r-r_0}{a}\right)^2, \qquad (3)$$

$$\Psi_1(r,\theta,\phi) = \Psi_1\left(\frac{r}{a}\right)^2 \cos(m\theta - n\phi). \tag{4}$$

$$T(r,\theta,\phi) = \Psi(r,\theta,\phi), \tag{5}$$

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Analytical test case (continued)

- boundary condition: $T(r = a) = \Psi(r = a)$
- ► analytical solution: parallel heat flux $\vec{q}_{||} = \hat{b}(\hat{b} \cdot \nabla T)$ vanishes as $(\nabla \Psi \times \hat{e}_{\phi}) \cdot (\nabla \Psi) = 0$
- reduces heat flux to $\vec{q} = \vec{q}_{\perp} = -n\chi_{\perp}\nabla T$
- ▶ heat diffusion equation $\nabla \cdot \vec{q} = -P \Rightarrow P = n\chi_{\perp}\nabla \cdot \nabla \Psi$ (to get the analytic solution)
- code runs with $\Psi_1/\Psi_0 = 10^{-2}$, m = 3, n = 2, $r_0/a = 1.2$
- resolution: 160 × 160 grid points

• error:
$$Err_{0,0} = \left[\frac{T_{num} - T_{analyt}}{T_{analyt}}\right]_{r=0.5a}^{0,0}$$

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Scale island width w_c

- competition between parallel and perp. transport
- ▶ for $w \approx w_c$, parallel \approx perpendicular transport
- ► scale island width $w_c = r_{res} \left(\frac{\chi_{||}}{\chi_{\perp}}\right)^{-1/4} \sqrt{\frac{8}{\epsilon_s s_s n}}$
 - n: toroidal mode number of perturbation
 - s_s: local magnetic shear
 - r_{res}: minor radius of resonant surface

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Finite difference scheme

Temperature grid (blue) and heat flux grid (green).



differential operators discretized with symm. 2nd order form
 e.g., the radial gradient of the temperature:

$$\left[\frac{\partial T}{\partial r}\right]_{i+\frac{1}{2},j+\frac{1}{2}} = \frac{T_{i+1,j+1} + T_{i+1,j} - T_{i,j+1} - T_{i,j}}{2\Delta r} + \mathcal{O}(\Delta r^2), \quad (6)$$

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